

# Luminosity, Flux and Magnitudes

Relativity and Astrophysics

Lecture 13

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## Outline

- Blackbody Radiation
- Flux and Luminosity
  - Inverse square law
- Magnitudes

## Radiation from objects

- All objects have internal energy which is manifested by the microscopic motions of particles.
- There is a **continuum** of energy levels associated with this motion.
- If the object is in **thermal equilibrium** then it can be characterized by a single quantity, it's **temperature**.
- An object in thermal equilibrium emits energy at all wavelengths.
  - resulting in a continuous spectrum
- We call this **thermal radiation**.

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## Blackbody Radiation

- A black object or **blackbody** absorbs all light which hits it.
- This blackbody also **emits** thermal radiation, e.g. photons!
  - Like a glowing poker just out of the fire.
- The amount of energy emitted (per unit area) depends only on the temperature of the blackbody.
- In 1900 Max Planck characterized the light coming from a blackbody.
- The equation that predicts the radiation of a blackbody at different temperatures is known as Planck's Law.

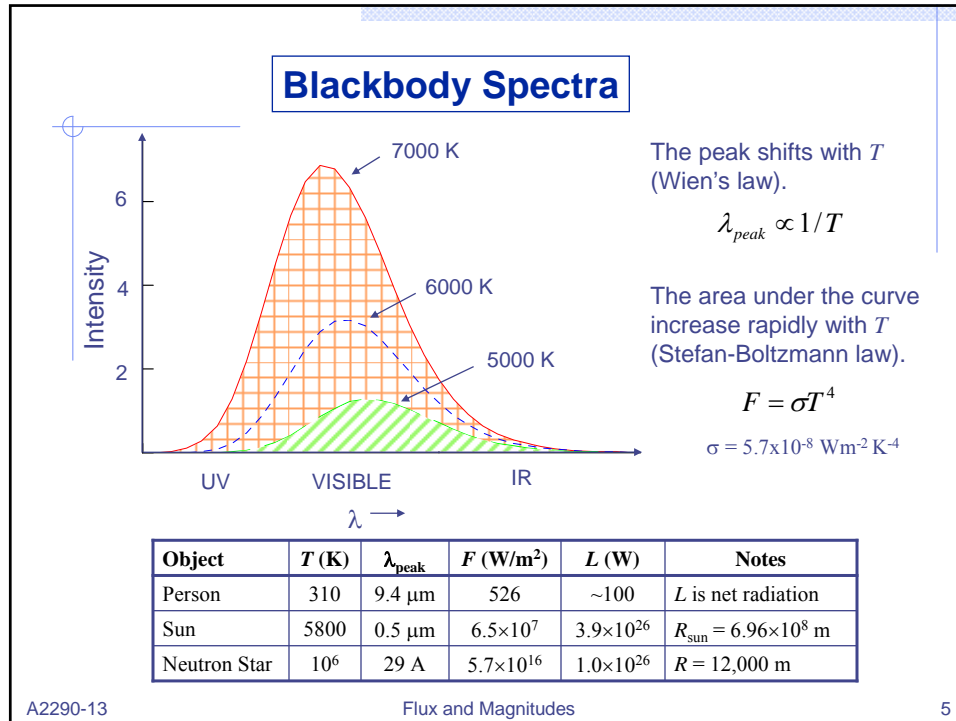
$$B_{\nu} = \frac{2h\nu^3}{c^2} \frac{1}{\exp(h\nu/kT)-1} \quad \text{W/m}^2/\text{Hz/sr}$$

- $h$  = Planck's constant,  $k$  = Boltzmann's constant,  $c$  = speed of light
- This is the power radiated per unit area in the frequency range  $\nu$  to  $\nu + d\nu$  into a unit solid angle ( $d\Omega = d\phi \sin\theta d\theta$  in spherical coordinates)

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## Properties of Blackbodies

- The peak emission from the blackbody moves to shorter wavelengths as the temperature increases (Wien's law).
 
$$\lambda_{peak} = 2900/T \quad \lambda \text{ in } \mu\text{m} \text{ and } T \text{ in K}$$
  - Hot objects look blue, cool ones look red
- The hotter the blackbody the more energy emitted per unit area at all wavelengths (Stefan-Boltzmann law).
 
$$F = \sigma T^4 \quad \sigma = 5.7 \times 10^{-8} \text{ Wm}^{-2} \text{ K}^{-4}$$
  - Note - bigger objects emit more radiation
- Except for their surfaces, stars behaves as a blackbodies

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## Energy Flux and Luminosity

- The Energy Flux,  $F$ , is the power per unit area radiated from an object.

$$F = \sigma T^4 \quad \text{W/m}^2 \text{ (at all } \lambda \text{)}$$

- The units are energy, area and time.
- Luminosity is the total energy radiated from star of radius  $R$  is given by:
- So the luminosity,  $L$ , is:

$$L = 4\pi R^2 \sigma T^4 \quad \text{Watts}$$

- If stars behave like blackbodies, stars with large luminosities must be very hot and/or very big.

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## Luminosity and Flux

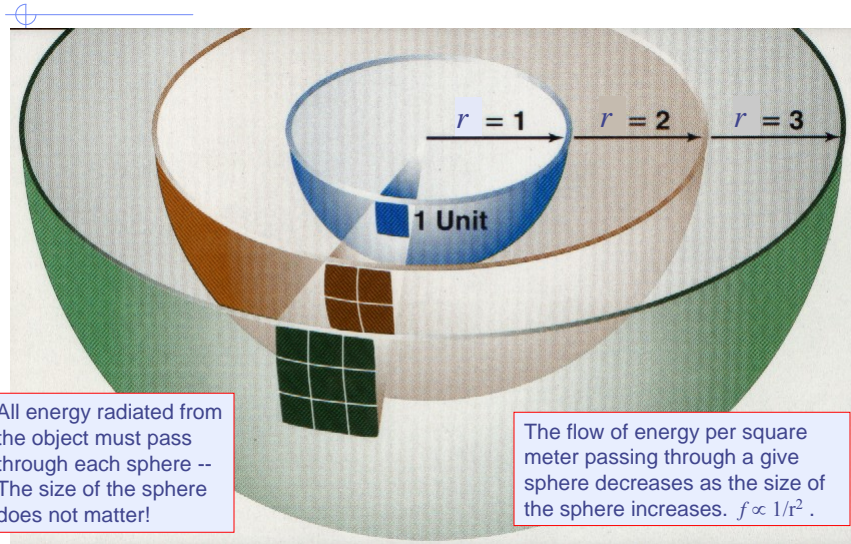
- Luminosity,  $L$ 
  - The total energy radiated from an object per second.
  - Measured in Watts
- Emitted Flux,  $F$ 
  - The flow of energy out of a surface.
  - Measured in Watts/m<sup>2</sup>
- Observed flux,  $f$ 
  - The power per unit area we receive from an object
  - Depends on the distance to the object.
  - Measured in W/m<sup>2</sup> e.g.  $f_{\text{sun}} = 1 \text{ kW/m}^2$
  - Also called flux or apparent brightness
- Meaning of Observed Flux
  - Make a sphere of radius,  $r$ , around an object (such as the Sun or a light bulb) which is radiating power.
  - All energy radiated from the object must pass through this sphere
    - ♦ The size of the sphere does not matter!

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## Flux falls off with distance



All energy radiated from the object must pass through each sphere -- The size of the sphere does not matter!

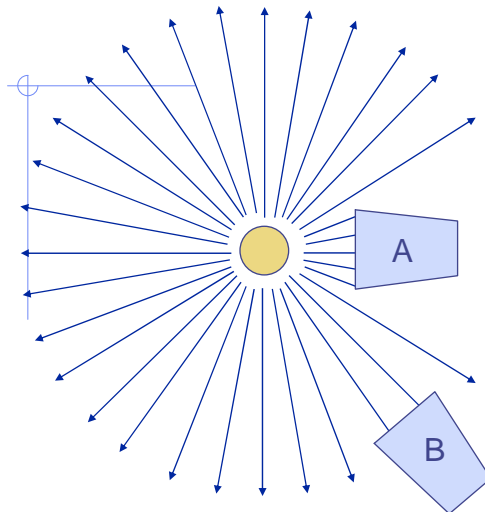
The flow of energy per square meter passing through a give sphere decreases as the size of the sphere increases.  $f \propto 1/r^2$ .

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## Flux: Sprinkler and Bucket Analogy



If bucket B is twice as far away as bucket A, it collects 1/4 as much water.

The sprinkler is the star, one of the buckets is the telescope (or your eye), and the water jets are the photons.

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## Inverse square law

- The flux,  $f$ , of energy through a sphere of radius,  $r$ , is given by

$$f = \frac{L}{4\pi r^2} \quad (\text{W/m}^2) \quad \text{Inverse square law}$$

where  $L$  is the luminosity of the object

- Why do we care about flux?
  - The flux is what we measure.
  - We use a telescope (or our eye) and measure a small fraction of the light passing through this sphere.

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## An illuminating example?

- A 100 W light bulb
  - about 1/5 of power goes into light
- It's total power output is always 100 W.
- It's apparent brightness to us depends upon how far away it is.
  
- For instance at 1 m the flux is:
  - Flux =  $0.08 \text{ W/m}^2$  [ $f = 100 \text{ W}/(4\pi(1\text{m})^2) \text{ W/m}^2$ ]
- If we double the distance away from the light bulb, the flux drops by a factor of 4.
  - At 2 m, the flux is  $0.02 \text{ W/m}^2$

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## Observed Flux – Distance Example:

- A star like the sun has an observed flux of  $2.4 \times 10^{-10} \text{ W/m}^2$ . If the flux of the sun at the Earth is  $1 \text{ kW/m}^2$ , how far away is the star?

$$L_{sun} = 4\pi d^2 f \Rightarrow L_{sun} = 4 \times 3.14 \times (1.5 \times 10^{11} \text{ m})^2 1000 \text{ W/m}^2$$

$$\Rightarrow L_{sun} = 3 \times 10^{26} \text{ W}$$

- Now

$$d = \sqrt{L_{sun} / 4\pi f} \Rightarrow d = \sqrt{3 \times 10^{26} \text{ W} / (4 \times 3.14 \times 2.4 \times 10^{-10} \text{ W/m}^2)}$$

$$\Rightarrow d = 3 \times 10^{17} \text{ m} = 10 \text{ pc}$$

- We could also have used ratios rather than compute  $L_{sun}$  first.

$$\frac{d_{star}}{d_{sun}} = \sqrt{\frac{f_{sun} L_{star}}{L_{sun} f_{star}}} = \sqrt{\frac{f_{sun}}{f_{star}}} \Rightarrow d_{star} = d_{sun} \sqrt{\frac{f_{sun}}{f_{star}}}$$

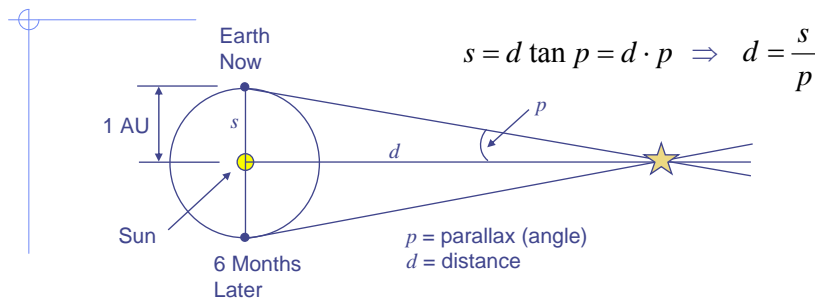
- Where we know everything but  $d_{star}$ .

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## Distances to stars: Stellar Parallax



- As stars get further away, their parallax becomes smaller.
- Parallax can not be measured to better than  $\sim 0.02''$  from the ground ( $d < 50 \text{ pc}$ ).
  - Interferometry is improving on this for selected applications
- Parallax is measured in arcseconds.
- Equations are for distances in AU and parsecs (pc), respectively

$$d(\text{AU}) = \frac{206265}{p(\text{arcsec})} \quad \text{or} \quad d(\text{pc}) = \frac{1}{p}$$

$1.0 \text{ arcsec} \Rightarrow 1 \text{ pc}$   
 $0.5 \text{ arcsec} \Rightarrow 2 \text{ pc}$

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## The closest stars

Star	Parallax (arcsec)	Distance (pc)	Luminosity ( $L_{\text{sun}}=1$ )
Proxima Centauri	0.763	1.31	$5 \times 10^{-5}$
$\alpha$ Centauri A	0.741	1.35	1.45
$\alpha$ Centauri B	0.741	1.35	0.4
Barnard's Star	0.522	1.81	$4 \times 10^{-4}$
Wolf 359	0.426	2.35	$2 \times 10^{-5}$
Lalande 21185	0.397	2.52	$5 \times 10^{-3}$
Sirius A	0.377	2.65	23
Sirius B	0.377	2.65	$2 \times 10^{-3}$

Parallax is motion of a star on the sky due to the Earth orbiting around the Sun. 1" parallax corresponds to 1 pc. The more distant a star the smaller the parallax. (1 pc = 3.26 lyr)

$$d(\text{pc}) = \frac{1}{p(\text{arcsec})}$$



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## Current status

- The Hipparcos Satellite (1989 – 1993)
    - Astrometry mission, produced two catalogs
  - Hipparcos catalog: ~120,000 stars
    - Measured parallaxes to better than 0.002"  $\Rightarrow d < 500$  pc
  - Tycho catalog: ~ 1,100,000 stars
    - Measured parallaxes and proper motions to ~ 0.025" (40 pc)
  - Tycho 2 catalog: 2,500,000 stars
    - Update version of Tycho catalog
    - Reprocessed raw Tycho data & used 144 other catalogs to obtain proper motions
    - Proper motions to 0.0025"/yr
- Parallaxes are the key to knowing distances in the universe.
  - Nearby stars are the stepping stone to measuring distance to everything else in the universe.
  - We can now compute the luminosity of stars!

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## Magnitudes

- We would like a way of specifying the relative brightness of stars
- Hipparchus
  - Devised a the magnitude system 2100 years ago to classify stars according to their apparent brightness.
  - He labeled 1080 stars as class 0, 1,.. 6.
  - 0 was the brightest, 1 the next brightest, etc.
- The magnitude scale is logarithmic.
- An **increase in magnitude** by 2.5 **means** an object is a factor of 10 **dimmer**, e.g.
  - a 0 mag star is 10 times brighter than a 2.5 mag star.
  - a 0 mag star is 100 times brighter than a 5 mag star.

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## Example magnitudes

Star	$m_v$	Designations
Sun	-26.8	
Sirius	-1.47	$\alpha$ CMa
Canopus	-0.72	$\alpha$ Car
Arcturus	-0.06	$\alpha$ Boo
Vega	0.03	$\alpha$ Lyr
Betegeuse	0.45	$\beta$ Ori
Altair	0.77	$\alpha$ Aqu
Deneb	1.26	$\alpha$ Cyg

- A dark adapted person with good eyesight can see to ~ 6<sup>th</sup> magnitude.
- Hubble Space Telescope can observed objects fainter than 30 mag.
  - $4 \times 10^9$  times fainter than the eye!

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## Fluxes and Magnitudes

- Flux is the power per unit area received from an object, e.g.  $f_{\text{sun}} = 1 \text{ kW/m}^2$
- If two stars, A and B, have fluxes,  $f_A$  and  $f_B$ , their magnitudes are related by

$$m_A - m_B = -2.5 \log(f_A / f_B)$$

- Thus if  $f_B / f_A = 10$ , then  $m_A - m_B = 2.5$
- We can also write the inverse relation

$$\frac{f_B}{f_A} = 10^{\frac{m_A - m_B}{2.5}}$$

- So that if  $m_A = 5$  and  $m_B = 0$ ,  $f_B / f_A = 100$ .

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## Absolute & Bolometric Magnitudes

- $m_v$  – apparent magnitude
  - How bright a star appears in the sky.
- $M_v$  – absolute magnitude
  - Brightness if the star were at 10 pc
  - This is an intrinsic property of the star!
- $M$  – absolute bolometric magnitude
  - Brightness at ALL wavelengths (and 10 pc).
- To get  $M_v$  or  $M$  we must know the distance to the star.
- Example:
  - Suppose a star has  $m_v = 7.0$  and is located 100 pc away.
  - It is 10 times the standard distance, thus, it would be 100 times brighter to us at the standard distance.
  - Or 5 magnitudes brighter  $\Rightarrow M_v = 2.0$

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## Example Absolute Magnitudes

Object	$m_V$	$M_V$
Sun:	-26.8	4.77
Full Moon:	-12.6	(32)
Sirius:	-1.47	1.4
Canopus:	-0.72	-3.1
Arcturus:	-0.06	-0.3
Deneb:	1.26	-7.2

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## The Distance Modulus Equation

- The relation between  $m_V$  and  $M_V$  is written in equation form as:

$$m_V - M_V = -5 + 5 \log_{10}(d) \quad (d \text{ in pc})$$

- $m_V - M_V$  is called the **distance modulus**.

- Examples:

- Deneb:  $m_V = 1.26$  and is 490 pc away.

$$m_V - M_V = -5 + 5 \log_{10}(d)$$

$$1.26 - M_V = -5 + 5 \log_{10}(490) = -8.5$$

$$\Rightarrow M_V = -7.2$$

- Sun:  $m_V = -26.8$ ,  $d = 1$  AU

$$-26.8 - M_V = -5 + 5 \log_{10}(1/206265)$$

$$\Rightarrow M_V = 4.8$$

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## Bonus: Deriving the S-B Law

- We get the Stefan-Boltzmann law by integrating the Planck function over all frequencies (area under the curve)

$$B = \int_0^{\infty} B_{\nu} d\nu = \int_0^{\infty} \frac{2h\nu^3}{c^2} \frac{d\nu}{\exp(h\nu/kT) - 1}$$

Let

$$x = \frac{h\nu}{kT} \Rightarrow \nu = \frac{kT}{h}x \quad \& \quad d\nu = \frac{kT}{h}dx$$

$$\Rightarrow B = \frac{2h}{c^2} \left( \frac{kT}{h} \right)^4 \int_0^{\infty} \frac{x^3 dx}{\exp(x) - 1} \Rightarrow B \propto T^4 \quad \text{W/m}^2/\text{sr}$$

- The total power emitted from a surface is proportional to the temperature to the fourth power – just as the S-B law
  - The constant of proportionality is not quite the S-B constant because we need to integrate over all solid angles to get it (which gives an additional factor of  $\pi$ ). The integral is related to the Riemann zeta function giving

$$\sigma = \pi \frac{2k^4}{c^2 h^3} \frac{\pi^4}{15} \Rightarrow \sigma = 5.67 \times 10^{-8} \quad \text{W/m}^2/\text{K}^4$$

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## Another form of Planck function

- Let's rewrite the Planck function in terms of power per unit wavelength rather than frequency interval

- Note that

$$B_{\nu} d\nu = B_{\lambda} d\lambda \quad \& \quad \nu = \frac{c}{\lambda} \Rightarrow d\nu = -\frac{c}{\lambda^2} d\lambda$$

- where  $B_{\lambda}$  is over the interval  $\lambda$  to  $\lambda + \Delta\lambda$  rather than  $\nu$  to  $\nu + \Delta\nu$ .
- This relationship must be true since the integral over frequencies and over wavelengths must be the same.

- Substituting gives

$$B_{\lambda} = B_{\nu} \frac{d\nu}{d\lambda} = \frac{2h\nu^3}{c^2} \frac{1}{\exp(h\nu/kT) - 1} \frac{c}{\lambda^2}$$

- Which yields for  $B_{\lambda}$

$$\Rightarrow B_{\lambda} = \frac{2hc^2}{\lambda^5} \frac{1}{\exp(hc/\lambda kT) - 1} \quad \text{W/m}^2/\text{m}/\text{sr} \text{ or } \text{W/m}^2/\mu\text{m}/\text{sr}$$

- Note that the units are now per m rather than per Hz.

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## Bonus: Deriving Wien's Law

- To find the peak we differentiate with respect to  $\lambda$  and setting this equal to zero to get the peak

$$\frac{dB_\lambda}{d\lambda} = \frac{2hc^2}{\lambda^5} \frac{1}{\exp(hc/\lambda kT) - 1} \left( -\frac{5}{\lambda} + \frac{\exp(hc/\lambda kT)}{\exp(hc/\lambda kT) - 1} \frac{hc}{\lambda^2 kT} \right) = 0$$

$$\Rightarrow \left( -5 + \frac{1}{1 - \exp(-hc/\lambda kT)} \frac{hc}{\lambda kT} \right) = 0$$

- Letting  $x = hc/\lambda kT$  and solving iteratively solving for  $x$ ,

$$x = 5(1 - \exp(-x)) \Rightarrow x = 4.965$$

- Putting back into the equation defining  $x$

$$\lambda T = \frac{hc}{xk} = \frac{6.626 \times 10^{-34} \text{ J-s} \times 3 \times 10^8 \text{ m/s}}{4.965 \times 1.38 \times 10^{-23} \text{ J/K}} \times \frac{10^6 \mu\text{m}}{\text{m}} \Rightarrow \boxed{\lambda T = 2900 \mu\text{m K}}$$

- which is Wien's law. Note that If we had solved for the peak in  $B_\nu$  (rather than  $B_\lambda$ ) this would yield,  $x = 2.82$  and the form of Wien's law is

$$\lambda T = 5100 \mu\text{m K} \quad B_\nu \text{ space}$$

- Both forms is correct since the peak is different between  $B_\nu$  and  $B_\lambda$  space.